



The Born-Oppenheimer Approximation:

A/c to the Born-Oppenheimer approximation the electronic motion in a molecule is assumed to take place independently of the nuclear motion.

Since the electronic and the nuclear motions are assumed to be separable we can write

$$\Psi(r, R) = \psi_e(r) \phi(R) \quad \text{--- (1)}$$

where $\psi_e(r)$ implies that the electronic wavefunction depends functionally on r and parametrically on R .

For simplicity we shall write ψ_e for $\psi_e(r)$ and ϕ_n for $\phi(R)$.

Let us define ψ_e as a function which satisfies the electronic Schrodinger equation

$$\hat{H}_e \psi_e = (\hat{T}_e + V_{en} + V_{ee}) \psi_e = E_e \psi_e \quad \text{--- (2)}$$

Since we know that the Schrodinger equation for the molecule is

$$\hat{H} \Psi(r, R) = (\hat{T}_e + \hat{T}_n + V_{en} + V_{ee} + V_{nn}) \Psi(r, R) = E \Psi(r, R) \quad \text{--- (3)}$$

Substituting $\Psi(r, R) = \psi_e \phi_n$ in eqⁿ (3) and writing the expressions for \hat{T}_e and \hat{T}_n explicitly we get,

$$\left(-\sum_i \frac{\nabla_i^2}{2} - \sum_A \frac{\nabla_A^2}{2m_A} + V_{en} + V_{ee} + V_{nn} \right) \psi_e \phi_n = E \psi_e \phi_n \quad \text{--- (4)}$$

$$\text{Now, } -\sum_i \frac{\nabla_i^2}{2} (\psi_e \phi_n) = -\phi_n \sum_i \frac{\nabla_i^2}{2} \psi_e \quad \text{--- (5)}$$

$$-\sum_A \frac{\nabla_A^2}{2m_A} (\psi_e \phi_n) = -\sum_A \frac{1}{2m_A} (\psi_e \nabla_A^2 \phi_n + \phi_n \nabla_A^2 \psi_e + 2 \nabla_A \psi_e \cdot \nabla_A \phi_n) \quad \text{--- (6)}$$

[since note that $d^2(xy) = xdy + ydx + 2dxdy$]

Now using eqⁿ (5) and (6) in eqⁿ (4) gives

$$-\phi_n \sum_i \frac{\nabla_i^2}{2} \psi_e - \left[\sum_A \frac{1}{2m_A} (\phi_n \nabla_A^2 \psi_e + 2 \nabla_A \psi_e \cdot \nabla_A \phi_n) \right]$$

$$- \psi_e \sum_A \frac{\nabla_A^2}{2m_A} \phi_n + (V_{en} + V_{ee} + V_{nn}) \psi_e \phi_n = E \psi_e \phi_n \quad \text{--- (7)}$$

The square-bracketed terms in eqⁿ (7) describes the coupling of electronic and nuclear motion. These terms are neglected in the BO approximation.